

THE FOURIER TRANSFORM ON $\mathcal{F}^*(\mathbb{R}^d)$

Suppose first that $T \in \mathcal{F}^*$ is a smooth & compactly supported function, $T \in C_c^\infty(\mathbb{R}^d)$. Then

$$\begin{aligned} \langle \hat{T}, \varphi \rangle &= \int_{\mathbb{R}^d} \underbrace{\beta^d}_{\mathbb{R}^d} \int_{\mathbb{R}^d} e^{-ix \cdot t} T(x) dx \varphi(t) dt = \boxed{\begin{array}{l} \text{The integrand is absolutely} \\ \text{summable, so we can invoke} \\ \text{Fubini to interchange the} \\ \text{integration order} \end{array}} \\ &= \int_{\mathbb{R}^d} T(x) \underbrace{\beta^d \int_{\mathbb{R}^d} e^{-ix \cdot t} \varphi(t) dt dx}_{= \hat{\varphi}(t)} = \langle T, \hat{\varphi} \rangle \end{aligned}$$

Since we've already proven that $\mathcal{F}: \mathcal{F} \rightarrow \mathcal{F}$ is continuous, we can now trivially define $\mathcal{F}: \mathcal{F}^* \rightarrow \mathcal{F}^*$ by duality.

Defn For $T \in \mathcal{F}^*(\mathbb{R}^d)$, define $\hat{T} = \widehat{\mathcal{F}T}$ by

$$\langle \hat{T}, \varphi \rangle = \langle T, \hat{\varphi} \rangle.$$

$$\text{Similarly, define } \check{T} = \widehat{\mathcal{F}^*T} \text{ by } \langle \check{T}, \varphi \rangle = \langle T, \check{\varphi} \rangle.$$

It follows immediately that

We proved earlier that if $\varphi \in \mathcal{F}$,
then $\mathcal{F}\mathcal{F}^*\varphi = \varphi$.

$$\langle \mathcal{F}^* \mathcal{F} T, \varphi \rangle = \langle \widehat{\mathcal{F}T}, \mathcal{F}^* \varphi \rangle = \langle T, \widehat{\mathcal{F}\mathcal{F}^*\varphi} \rangle = \langle T, \varphi \rangle$$

So \mathcal{F}^* is the inverse of \mathcal{F} on \mathcal{F}^* . To sum up:

Propn The map $\mathcal{F}: \mathcal{F} \rightarrow \mathcal{F}$ is continuous, bijective,
and ~~so~~ its inverse is continuous as well.

Example $\widehat{\delta} = S \leftarrow$ Dirac delta function

$$\langle \hat{T}, \varphi \rangle = \langle T, \hat{\varphi} \rangle = \hat{\varphi}(0) = \beta^d \int \varphi e^{-i0 \cdot x} \varphi(x) dx = \beta^d \int \varphi(x) dx = \langle \beta^d, \varphi \rangle$$

$$\text{So } \widehat{\delta} = \frac{1}{(2\pi)^d} \delta/2$$

The Fourier transform on $L^2(\mathbb{R}^d)$

For now, we switch to complex-valued functions on \mathcal{F} & \mathcal{F}^* .

Nothing that we've done changes.

Def' Let (\cdot, \cdot) denote the ~~L^2~~ inner product, $(u, v) = \int_{\mathbb{R}^d} \overline{u(x)} v(x) dx$. Define $\tilde{L}^2(\mathbb{R}^d)$ by taking the closure of \mathcal{F} w.r.t. the $\|\cdot\|_{\tilde{L}^2(\mathbb{R}^d)}$ norm.

Lemma If $\varphi, \psi \in \mathcal{F}(\mathbb{R}^d)$, then $(\varphi, \psi) = (\hat{\varphi}, \hat{\psi})$.

Proof $(\varphi, \psi) = \int \bar{\varphi} \psi = \langle \bar{\varphi}, \psi \rangle = \langle \bar{\varphi}, \check{\psi} \rangle = \langle \frac{\vee}{\bar{\varphi}}, \check{\psi} \rangle \iff$

$$\text{We have } \check{\bar{\varphi}}(t) = \beta^d \int e^{ixt} \bar{\varphi}(x) dx = \overline{\beta^d \int e^{-ixt} \varphi(x) dx} = \overline{\hat{\varphi}(t)} \text{ so}$$

$$(\varphi, \psi) = \langle \check{\bar{\varphi}}, \check{\psi} \rangle = \langle \hat{\bar{\varphi}}, \hat{\psi} \rangle = (\hat{\varphi}, \hat{\psi}).$$

As a consequence of the lemma: $\|F\varphi\|_2^2 = (\hat{\varphi}, \hat{\varphi}) = (\varphi, \varphi) = \|\varphi\|_2^2$

Thus $F: \mathcal{F} \rightarrow \tilde{L}^2(\mathbb{R}^d)$ is an isometric map.

Since \mathcal{F} is dense in $L^2(\mathbb{R}^d)$, we can ~~extend~~ ^{uniquely} F to all of $L^2(\mathbb{R}^d)$.

Def' We define $F: L^2(\mathbb{R}^d) \rightarrow \tilde{L}^2(\mathbb{R}^d)$ as follows:

For $f \in L^2(\mathbb{R}^d)$, pick $\varphi_n \in \mathcal{F}$ s.t. $\|F\varphi_n\|_2 \rightarrow 0$. Set $\hat{f} = \lim_{n \rightarrow \infty} \hat{\varphi}_n$.

Define $F^*: \tilde{L}^2 \rightarrow L^2$ analogously, $\check{f} = \lim_{n \rightarrow \infty} \check{\varphi}_n$.

\uparrow Limit is in L^2 -sense.

It follows immediately that $F^* = F^{-1}$ and thus that F is bijective.
To summarize:

Thm The Fourier transform $F: L^2(\mathbb{R}^d) \rightarrow L^2(\mathbb{R}^d)$ is

- * Bijective (one-to-one & onto)
- * Preserves the norm: $\int |u(x)|^2 = \int |\hat{u}(x)|^2$
- * Preserves the inner product: $\int \overline{u(x)}v(x) = \int \overline{\hat{u}(t)}\hat{v}(t) dt$
- * Is continuously invertible.

In short, it is a "unitary map" or a "Hilbert space isomorphism".

Note We defined \hat{f} as the limit $\hat{f} = \lim \hat{\phi}_n$, where $\phi_n \in S$.

In practice, it is enough to pick $\phi_n \in L^2 \cap L^1$, for instance

$$\hat{f}(t) = \lim_{R \rightarrow \infty} \int_{|x| \leq R} e^{-ix \cdot t} f(x) dx = \lim_{\epsilon \rightarrow 0} \int e^{-ix \cdot t} e^{-\epsilon|x|^2} f(x) dx$$

What is the spectrum of F ?

First we note that since F is unitary, we must have $\sigma(F) \subseteq \{\lambda \in \mathbb{C} : |\lambda|=1\}$.

Moreover, ~~note that~~ $F^2 = FF^*F = F^*F = I$, and so $F^4 = I^2 = I$.

This ~~implies~~ implies that if $\lambda \in \sigma(F)$, then $\lambda^4=1$ and so $\lambda \in \{1, -1, i, -i\}$.

(Formally: $F^4 = (\int_{\sigma(F)} \lambda dP(\lambda))^4 = \int_{\sigma(F)} \lambda^4 dP(\lambda) = I \Rightarrow \lambda^4=1$)

In fact the four numbers $1, -1, i, -i$ are all evals of infinity multiplicity and we have explicit formulas for the eigenvectors!

In one dimension, those are the so called Hermite functions $(\phi_n)_{n=0}^\infty$, which form an ON-basis for $L^2(\mathbb{R})$.
 $\phi_n(x) = \alpha_n e^{x^2/2} \left(\frac{d}{dx}\right)^n e^{-x^2} = \beta_n H_n(x) e^{-x^2/2}$ where H_n are the Hermite polynomials.

(The H_n are constructed for instance by applying Gram-Schmidt to $\{1, x, x^2\}$ w.r.t. $\langle u, v \rangle = \int \bar{u} v e^{-x^2} dx$)

It turns out that $\mathcal{F}\varphi_n = (-i)^n \varphi_n$ and so

$$\begin{aligned} L^2(\mathbb{R}^d) &= \ker(\mathcal{F} - I) \oplus \ker(\mathcal{F} + iI) \oplus \ker(\mathcal{F} - iI) \oplus \ker(\mathcal{F} + 2iI) = \\ &= \text{span}(\varphi_{4n}) \oplus \text{span}(\varphi_{4n+1}) \oplus \text{span}(\varphi_{4n+2}) \oplus \text{span}(\varphi_{4n+3}) \end{aligned}$$

With this knowledge, we could also have defined \mathcal{F} via

$$\mathcal{F}\left[\sum_{n=0}^{\infty} \alpha_n \varphi_n\right] = \sum_{n=0}^{\infty} \alpha_n (-i)^n \varphi_n.$$

Sobolev Spaces

When is $\partial^\alpha \mathcal{F}$ an L^2 function?

$$\partial^\alpha \mathcal{F} \in L^2 \Leftrightarrow \int |\partial^\alpha \mathcal{F}|^2 < \infty \Leftrightarrow \int |(-it)^{\alpha} \hat{f}(t)|^2 dt < \infty \Leftrightarrow (-it)^{\alpha} \hat{f} \in L^2$$

Let s be a non-negative integer and define the Sobolev space $H^s(\mathbb{R}^d)$ by

$$\begin{aligned} H^s(\mathbb{R}^d) &= \{f : \partial^\alpha f \in L^2 \quad \forall \alpha : |\alpha| \leq s\} = \\ &= \{f : (-it)^s \hat{f} \in L^2 \quad \forall \alpha : |\alpha| \leq s\} = \\ &= \{f : (1+t^2)^{s/2} \hat{f} \in L^2\} = \{f : \int (1+t^2) |\hat{f}(t)|^2 dt < \infty\}. \end{aligned}$$

The definition is readily extended to any $s \in \mathbb{R}$, not just positive integers.

Lemma If $f \in H^s(\mathbb{R}^d)$, for some $s > d/2$, then $f \in C_0(\mathbb{R}^d)$.

More generally, if $f \in H^s(\mathbb{R}^d)$ for some $s > \frac{d}{2} + k$, then $f \in C^k_0(\mathbb{R}^d)$.