The Last Invariant Torus J.D. Meiss University of Colorado

Symplectic Twist Maps

Poincaré Sections of Natural Hamiltonians

 $\frac{dz}{dt} = J\nabla H(z) \quad z = (q, p) \qquad J = \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix}$

H(q,p) = T(p) + V(q)

Symplectic Twist Maps

Froeshlé Map

 $x' = x + Ty' \mod 1$ $y' = y - \nabla V(x)$

Symplectic if T symmetric $f^*\omega = \omega$ Exact Symplectic (zero net flux)

$$\omega = dx \wedge dy = \sum_{i=1}^{n} dx_i \wedge dy_i$$

$$f^*(ydx) - ydx = dL$$

Twist Condition: $T = K^{-1} > 0$ $L(x, x') = \frac{1}{2}(x' - x)^T K(x' - x) - V(x)$ $y = -L_1 = K(x' - x)$ $y' = L_2 = K(x' - x) - \nabla V(x)$

T = I $V(x) = \frac{1}{4\pi^2} \left(a \cos(2\pi x_1) + b \cos(2\pi x_2) + c \cos(2\pi (x_1 + x_2)) \right)$



y₂ projection

 x_1 projection

a=0.1, b=0.2 c = 0.1



Frequency Maps

Frequency: Action to $\omega \quad \Omega : \mathbb{R}^n \to \mathbb{R}^n$ Froeshlé case $\Omega(y) = Ty$ KAM: assumes Ω is a local diffeomorphism No good theory of multi-dim cont. fractions(100) Best approximates? Periodicity? (210)Farey Tree generalization LR binary encoding of frequency vectors (101) L (110)LL $(1,\sigma^2,\sigma)$ (112) natural self-similar patterns R

- spiral mean $\sigma^3 = \sigma + 1$
- Diophantine vector (σ , σ^2 , 1)

Kim, S. and S. Ostlund (1986). "Simultaneous Rational Approximations in the Study of Dynamical Systems." Phys. Rev. A 34: 3426-3434.

(001)

(120)

(010)

RR

(001)

A Tale of Three Methods

Converse KAM Theory Frequency Analysis Crossing Time

#1 Converse KAM Theory

Nonexistence criteria

- * Birkhoff's Theorem: every rotational invariant circle of an area-preserving twist map is a graph $\{(x, S'(x)) : x \in \mathbb{T}^2\}$
 - Confinement Corollary: if all orbits below y = a remain below y = b, then there is a rotational invariant circle in (a,b).
 - By twist condition, S is Lipschitz

Cone Criterion
$$Df = \begin{pmatrix} 1 - F'(x) & 1 \\ -F'(x) & 1 \end{pmatrix}$$

 $v_{+} = Df \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$
 $v_{-} = Df^{-1} \begin{pmatrix} 0 \\ -1 \end{pmatrix} = \begin{pmatrix} 1 \\ -1 + F' \end{pmatrix}$
 $v'_{+} = Dfv_{+} = \begin{pmatrix} 2 - F' \\ 1 - F' \end{pmatrix}$

Converse KAM Theory

Mather: no rotational circles if k > 4/3

+ Rigorous numerics: if k > 63/64

Mather, J. N. (1984). "Non-Existence of Invariant Circles." <u>Ergodic Theory and Dynamical Systems</u> 4: 301-309. MacKay, R. S. and I. C. Percival (1985). "Converse KAM: Theory and Practice." <u>Comm. Math. Phys.</u> 98: 469-512.

Higher Dimensions?

* Analogue of Birkhoff's theorem: Lagrangian invariant tori on which dynamics is chain recurrent are graphs $\{(x, \nabla S(x)) : x \in \mathbb{T}^2\}$

Bialy, M. L. and L. Polterovich (1992). "Hamiltonian Systems, Lagrangian Tori and Birkhoff's Theorem." Math. Ann. 292: 619-627.

- Every orbit on an invariant Lagrangian graph as minimal action
- + Simplest criterion: need |a| + |c| < 2 and |b| + |c| < 2 for tori.

Higher Dimensions Froeshlé Map, a = 0.05, b = 0.2, c = 0.02



MacKay, R. S., J. D. Meiss and J. Stark (1989). "Converse KAM Theory for Symplectic Twist Maps." <u>Nonlinearity</u> 2: 555-570.

#2 Frequency Analysis

KAM Tori have well-defined rotation vectors

 $\omega = \lim_{t \to \infty} \frac{x_t - x_0}{t}$

Each torus crosses every plane x = const, so sufficient to look on one

for example, symmetry plane Fix(S₁) ={x = 0}.
Finite time approximate frequencies
May use windowed-FFT methods (Laskar)

Frequency Analysis

Birkhoff: invariant circles are graphs

Twist \Rightarrow frequency is monotone on circles



Fig. 2(a)-(d). Variation of the fundamental frequency ν for the standard mapping (13) for different values of the parameter *a*, in the vicinity of the golden rotation number ν_0 which corresponds to the zero dotted line. The origin in the *x* scale is arbitrarily taken to be $x_0 = 4.17655$. The origin of frequencies is the golden value $\nu_g = \frac{1}{3}(3 - \sqrt{5})$. The unit for ν and *x* is 10⁻⁶. If $x_1 < x_2$ and $\nu(x_1) > \nu(x_2)$, we can conclude that there exist no KAM invariant curves of irrational rotation number between $\nu(x_2)$ and $\nu(x_1)$. In fig. 9b, we can see that the golden invariant curve does not persist for a = 0.9718.

Laskar, J. (1993). "Frequency Analysis for Multi-Dimensional Systems. Global Dynamics and Diffusion." Physica D 67: 257-283.

Frequency Analysis

Birkhoff: invariant circles are graphs

Twist \Rightarrow frequency is monotone on circles



Fig. 2(a)-(d). Variation of the fundamental frequency ν for the standard mapping (13) for different values of the parameter *a*, in the vicinity of the golden rotation number ν_0 which corresponds to the zero dotted line. The origin in the *x* scale is arbitrarily taken to be $x_0 = 4.17655$. The origin of frequencies is the golden value $\nu_g = \frac{1}{3}(3 - \sqrt{5})$. The unit for ν and *x* is 10⁻⁶. If $x_1 < x_2$ and $\nu(x_1) > \nu(x_2)$, we can conclude that there exist no KAM invariant curves of irrational rotation number between $\nu(x_2)$ and $\nu(x_1)$. In fig. 9b, we can see that the golden invariant curve does not persist for a = 0.9718.

Laskar, J. (1993). "Frequency Analysis for Multi-Dimensional Systems. Global Dynamics and Diffusion." Physica D 67: 257-283.

Higher Dimensions? Froeshlé Map, a = b = 1.3



Laskar, J. (1993). "Frequency Analysis for Multi-Dimensional Systems." <u>Physica D</u> 67: 257-283. Dullin, H. R. and J. D. Meiss (2003). "Twist Singularities for Symplectic Maps." <u>Chaos</u> 13: 1-16.

#3 Crossing Time





 $N = \frac{C}{(k - k_c)^{\mu}}$

Chirikov: fit to data $k_c = 0.989, \mu = 2.55$

MMP : cantorus flux $k_c = 0.971635,$ $\mu = \eta = 3.01177$ C = 25

Chirikov, B. V. (1979). "A Universal Instability of Many-Dimensional Oscillator Systems." <u>Phys. Rep.</u> 52: 265-379.
MacKay, R. S., J. D. Meiss and I. C. Percival (1984). "Transport in Hamiltonian Systems." Physica D 13: 55-81.

Higher Dimensions?



Volume Preserving Maps

Magnetic Field line flows $\frac{dx}{dt} = B(x,t)$ $\nabla \cdot B = 0$ Incompressible Fluids $\frac{dx}{dt} = v(x,t)$ $\nabla \cdot v = 0$

Poincaré Map for Periodic Time dependence: V.P.

Invariant Tori

KAM theory applies to one-action, *n*-angle, exact V.P.Maps $(x,z) \in \mathbb{T}^n \times \mathbb{R}$ Cheng, C.-Q. and Y.-S. Sun (1990). "Existence of Invariant Tori in Three
Dimensional Measure-Preserving Mappings." Celestial Mech. 47(3):
275-292.

$$x' = x + \Omega(z') + \varepsilon g_1(x, z) \mod 1$$
$$z' = z + \varepsilon g_2(x, z) \qquad \qquad \int_{\mathbb{T}^n} g_2 dx = 0$$

providing det $(D\Omega, D^2\Omega) \neq 0$.

Cantor sets of invariant tori for $|\varepsilon| << 1$, though cannot be identified by fixed frequency vector

 $\omega \in \mathbb{R}^n$

A Residue Criterion?

KAM theory applies (Cheng & Sun)

- However, can't fix the frequencies!
- Is there a last torus? Self-similarity?
 - What rotation vector plays the role of the golden mean? Perhaps a cubic irrational: spiral mean $\sigma^3 = \sigma + 1$?
- Are there cantori?

Some results: Anti-integrability by Li & Malkin
 Is there an algebraic singularity in the crossing time?

Standard VP Map

3D, one-action map

$$\begin{aligned} x_1' &= x_1 + \Omega_1(z') , \\ x_2' &= x_2 + \Omega_2(z') , \\ z' &= z + \varepsilon g(x) , \end{aligned}$$

For "twist condition" need nonzero curvature

$$\Omega(z) = (z + \gamma, -\delta + \beta z^2) .$$

Dullin, H. R. and J. D. Meiss (2010). "Resonances and Twist in Volume-Preserving Mappings." <u>Disc. Cont. Dyn. Sys.</u> submitted.

Forcing any generic periodic function

 $g(x) = -a\sin(2\pi x_1) - b\sin(2\pi x_2) - c\sin(2\pi(x_1 - x_2))$

 $\Omega:\mathbb{R}\to\mathbb{R}^2$

0.

0.

ω

$$\Omega(z) = (z + \gamma, -\delta + \beta z^2)$$

Resonances

 $m \cdot \omega = n$

Diophantine Condition

$$|m \cdot \omega - n| > \frac{C}{|m|^{\tau}} \qquad \text{-0.2}$$

 $\mathcal{R} \equiv \left\{ \omega \in \mathbb{R}^d : m \cdot \omega = n \text{ for some } (m, n) \in \mathbb{Z}^{d+1} \setminus \{0\} \right\}$

-0.

Resonances Driven (1,0,0), (0,1,0) and (1,1,0) resonances $\gamma = \frac{1}{2}(\sqrt{5}-1) \approx 0.61803$ $\beta = 2$ a = b = c = 1.0 $\delta = 0.1$





Frequency Analysis

Simplest numerical estimate

$$\omega_T(x_0, z_0) = \frac{x_T - x_0}{T}$$



 $x_0 = (0,0)$ $T = 10^5$

Frequency Maps $\varepsilon = 0.01$



The Last Torus

Easiest case: z-periodic structure

 $\Omega(z+1) = \Omega(z) + m$

Any invariant set for $z \in [0,1]$ repeated in [k,k+1]. To test for invariant tori need to bound vertical extent

 $\Delta(\mathcal{C}) = max_{\mathcal{C}}(z) - min_{\mathcal{C}}(z)$ Experiments indicate Δ_{max} is small, say < 0.1.

: if there are no tori in $[0, 1 + \Delta_{max}]$, there are none.

Periodicity in z

Periodic structure in vertical: $\Omega(z+1) = \Omega(z)+m$



Periodicity in z

Periodic structure in vertical: $\Omega(z+1) = \Omega(z)+m$















Crossing Time $t_c = \min\{t > 0 : |z_t - z_0| \ge 1\}$

 $\delta = 0.1$

 $\delta = 0.3$



10 initial conditions for each ε

The Last Torus?



 $\delta = 0.1$ $\epsilon = 0.02725$

 $z_0 = -0.0560$

 $\omega \approx (0.618681, -0.085983) \\ = ([0, 1^7, 5, 1^2,], [-1, 1^2, 10, 1^3, 2^2,])$

The Last Torus?

Analog of Greene's Self-similarity?

Analog of the golden mean?

Analog of Birkhoff's 2nd theorem?

- Rotational tori need not be graphs
- Can one explicitly bound their vertical extent?

Flux scaling exponent?

Analog of Aubry-Mather: Cantori?

Thanks!

Differential Dynamical Systems

James D. Meiss







and buy my book!